



M
full phase space

$\bar{M}$
Constraint
Surface
 $C(x) = E(x) = 0$

$\hat{M}$
reduced phase
space

= gauge equivalence
of points on $\bar{M}$ classes

Definition

$F(q, p)$

is said to be a

Strong Dirac observable $\Leftrightarrow \{C(x), F\} = \{C_+(x), F\} = 0$
 $\forall x$

weak

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"

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on $\overline{M}$

Quantisation of Systems with first class constraints

- Suppose we have first class
constraints C_I

- Suppose that $\{C_I, C_J\} = f_{IJ}^K C_K$

Structure
constants

(indep of $T, \bar{T}$)

$\{ \} = 0$

- Choose some ^{sub-Poisson} algebra $\mathcal{A}$ of "elementary" variables also closed under complex conjugation

$\mathcal{A}$ should separate the points on the phase space

- Impose CCR's

$$[\hat{a}, \hat{b}] = i\hbar \{a, b\}$$

Impose $\ast$ relations

$\Rightarrow$

$\hat{\mathcal{A}}$

quantum algebra

$$\hat{a} + \hat{a}^\dagger = \hat{a}$$

Find Hilbert space representation of $\hat{A}$
 such that $\hat{C}_I$ can be defined as
 s.a. operator on $\mathcal{H}$ (Kinematical HS)
 and s.t. $[\hat{C}_I, \hat{C}_J] = i\hbar f_{IJ} \hat{C}_K$

One would like to select physical states
 weak $\hat{C}_I \psi = 0$, $\psi \in \mathcal{H}$ (quantum analog
 of $C_I = 0$)

- look for solutions as follows

$$(C_I \Psi = 0 \quad \|\Psi\|_H = \infty)$$

(class) constraints

Rigging Map

Find dense, invariant subspace

$$\overline{\Phi} \text{ of } H \subset I \quad \hat{C}_I \overline{\Phi} \subset \overline{\Phi}$$

Let $\overline{\Phi}^*$ be the algebraic

dual of $\overline{\Phi}$. $\Psi \in \overline{\Phi}^*$ distribution
on $\overline{\Phi}$, $\Psi[f]$

$$\phi \subset \mathcal{H} \subset \phi^*$$

$$\Psi(f) = \langle \Psi, f \rangle$$

$$\bar{\phi} \rightarrow \bar{\phi}^*$$

case by case Ansatz

$$f \mapsto \int_G d\mu(g) \langle U(g) f, \cdot \rangle$$

$$U(g) = \exp \left(i \sum_I g^I \hat{C}_I \right)$$

unitary operators

μ Haar measure on the group G generated by the Lie algebra generators C_I

$$d\mu(g g_0) = d\mu(g)$$

Want to solve
of

$$\hat{C}_I \psi = 0$$

$$U(g) \psi = \psi$$

$$\psi(U(g)f) = \psi(f)$$

instead

Choose some generators

τ_I for Lie Algebra

define $g(t)$

$$= \exp(t \tau_I)$$

$$d\mu(g) = d\tilde{\mu}(g(t))$$

$$\Psi(U(g)f) \quad \text{follows}$$

$$= \int d\mu(g') \langle U(g')\Psi, U(g)f \rangle$$

$$= \int d\mu(g') \langle \underbrace{U(g^{-1})U(g')}_{U(g^{-1}g')} \Psi, f \rangle$$

$$= \int d\mu(gg') \langle U(g')\Psi, f \rangle$$

$$= \Psi(f)$$

sub-Poisson

$$\langle \eta(f), \eta(f') \rangle_{\text{phys}}$$

$$\langle \psi, \psi' \rangle$$

$$\eta = \int d\mu(g) \langle U(g) f, f' \rangle_{\mathcal{H}}$$

case
case
Example

$$q_1, p_1, q_2, p_2 \quad \mathcal{L} = p_2^2, \quad \mathcal{H} = L_c(\mathbb{R}^2, dx, dx_2)$$

$$\left(\hat{p}_2 \psi \right) (x_1, x_2) = i \hbar \frac{\partial \psi}{\partial x_2} (x_1, x_2) = 0 \Rightarrow \psi = \psi(x_1)$$

Lie algebra generators $\mathcal{L}_1$

However $\|\psi\|^2 = \int dx_1 |\psi(x_1)|^2 \left(\int dx_2 \right)$

Let $\phi = \mathcal{S}(\mathbb{R}^2)$ smooth functions of rapid decrease at ∞

$\phi' = \mathcal{S}'(\mathbb{R}^2)$

Schwarz space of tempered distribution

$\uparrow$ cont. fns on $\mathbb{R}$
 $\phi * \psi$ algebraic dual

$$\eta(t) = \int \frac{dt'}{2\pi} \langle e^{it\hat{p}_2} f \rangle$$

follows

$$= \langle \delta(\hat{p}_2) f \rangle$$

$$\hat{p}_2 \delta(\hat{p}_2) = 0$$

$$\eta(t) [e^{it\hat{p}_2} f']$$

$$= \int \frac{dt'}{2\pi} \langle e^{i(t'-t)\hat{p}_2} f, f' \rangle$$

$$= \int \frac{dE}{2\pi} \langle e^{it'} f, f' \rangle$$

Sub - Poisson

$$\langle \gamma(t), \gamma(t') \rangle_{\text{phys}} = \int \frac{dt}{2\pi} \langle e^{it\vec{p}_2} f, f' \rangle_{L_2(d\vec{x})}$$

$$= \int dt \int d^2x \frac{f(x_1, x_2 - t) f'(x_1, x_2)}{2\pi}$$

$$\Rightarrow \int dx_1 \left[\int dx_2 f(x_1, x_2) \right] \left[\int dx'_2 f'(x_1, x'_2) \right]$$

$$\Rightarrow \gamma(t) \cong \int dx_2 f(x_1, x_2) \quad H = L_c(\mathbb{R}^2 / dx_1, dx_2)$$

$$H_{\text{phys}} \cong L_2(\mathbb{R}, dx_1)$$

Lie algebra generators L_J